

# Properties of Heavy Flavor Hadrons in ATLAS

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*HADRON 2007, Frascati*

*12 October 2007*

The Large Hadron Collider (7 TeV - on - 7 TeV proton-proton collisions) will provide unprecedented energy scales for QCD study.

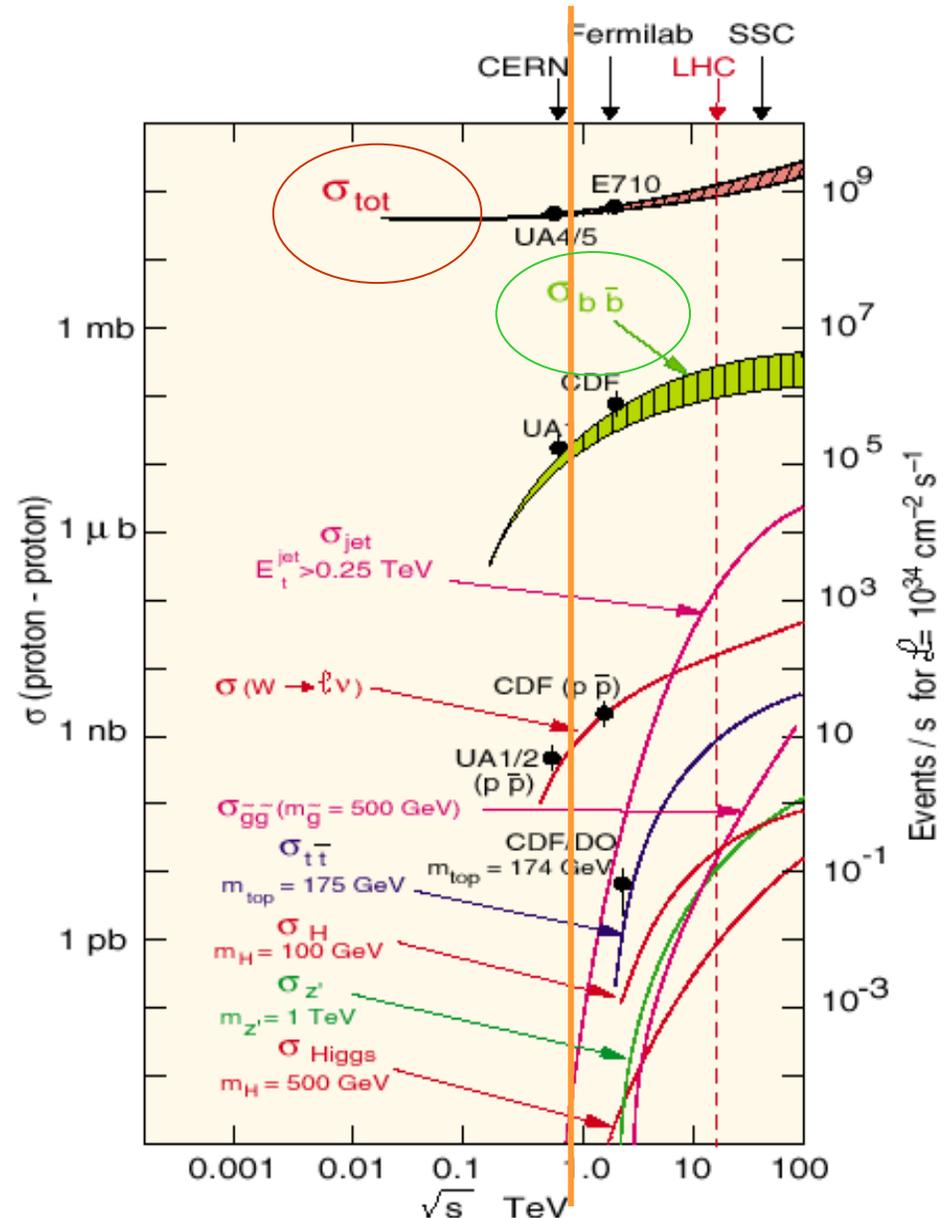


LHC AERIAL VIEW

ATLAS, one of 4 LHC detectors, will cover collisions to pseudorapidity  $|\eta| \leq 2.4$  for high precision measurements.

# ATLAS, the LHC, and B Physics

- LHC has a 27 km circumference, 40 kHz crossing rate
- The total  $b\bar{b}$  production cross section is 500  $\mu\text{b}$ : 1  $b\bar{b}$  pair in every 100 collisions.
- Luminosity expectation: 10  $\text{fb}^{-1}$  per year (@  $L=10^{33}/\text{cm}^2/\text{s}$ ) in Years 1-3, 100  $\text{fb}^{-1}$  per year subsequently.



# The ATLAS Detector

Tracking ( $|\eta| < 2.5$ ,  $B = 2T$ ) :

Silicon pixels and strips

Transition Radiation Detector (tracking and  $e/\pi$  separation)

Calorimetry ( $|\eta| < 5$ ) :

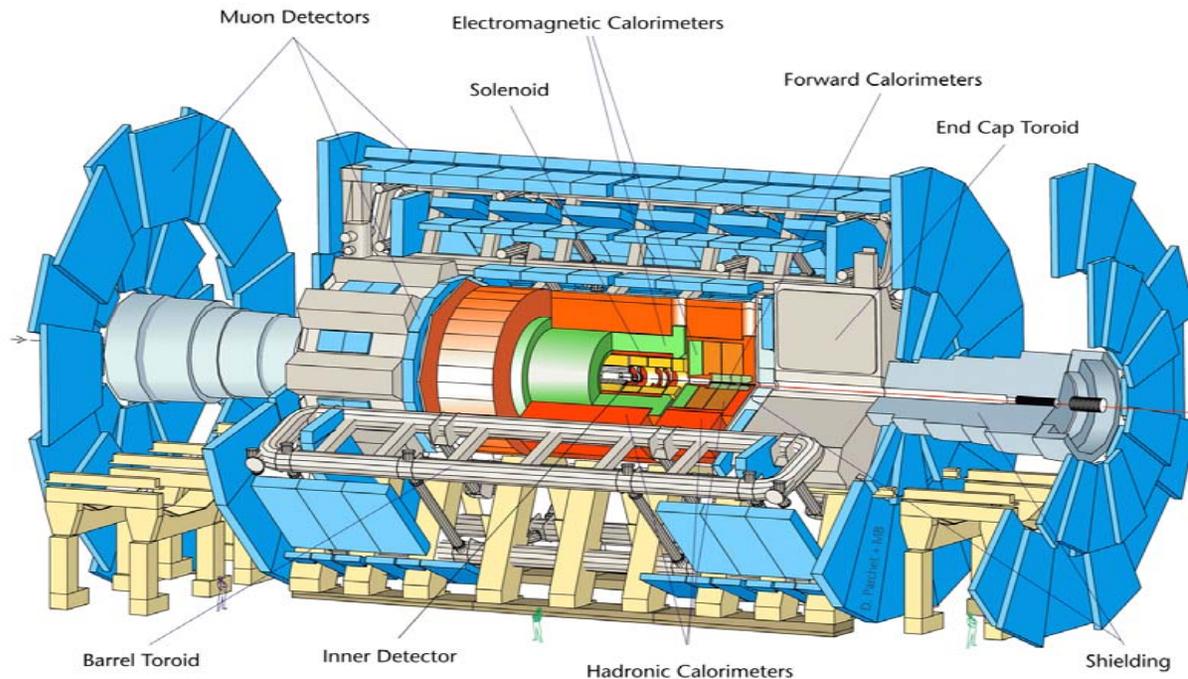
EM : Pb-LAr

HAD: Fe/scintillator (central), Cu/W-LAr (fwd)

Muon Spectrometer ( $|\eta| < 2.7$ ) :

Air-core toroids with muon chambers

- 46 m long
- 22 m diameter
- 7000 t total weight
- 2T solenoid and 0.5 T toroid
- $10^8$  electronics channels
- 3000 km of cables.



# Inner Detector

## Pixels:

- $(0.8 \times 10^8 \text{ channels})$
- $\sigma_\phi = 12 \mu\text{m}$ ,  $\sigma_z = 66 \mu\text{m}$

## Silicon Tracker (SCT):

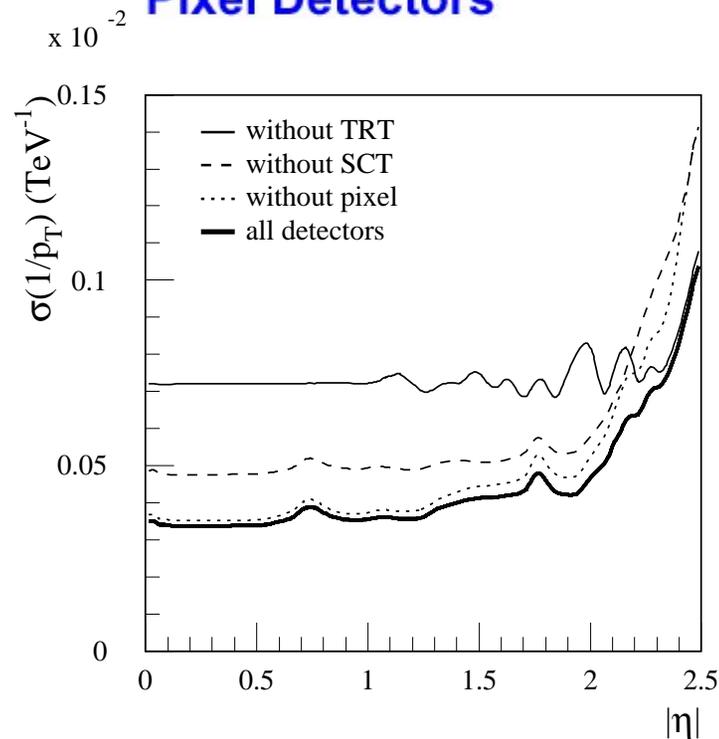
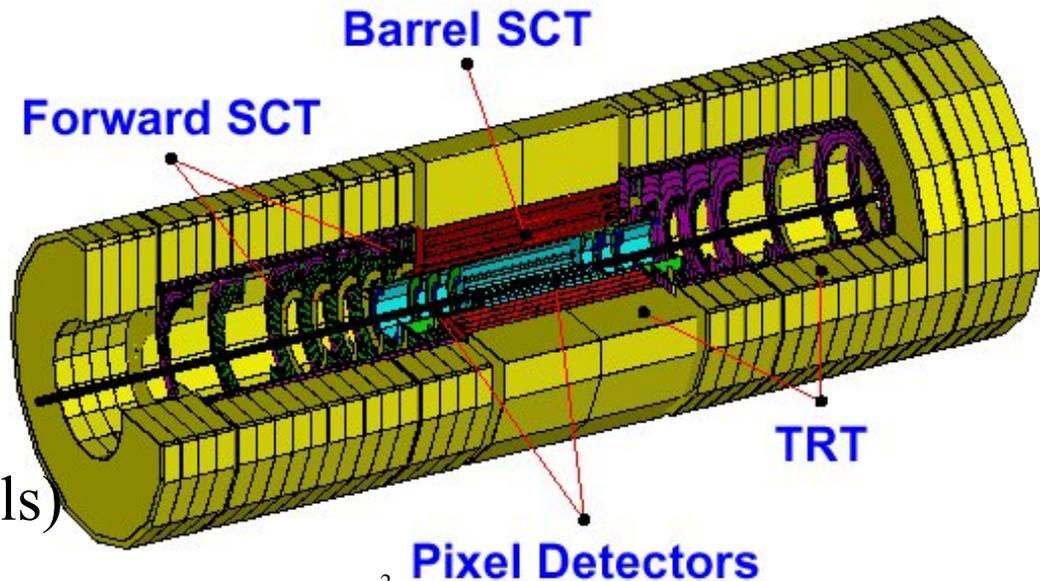
- $5\text{cm} < \text{radii} < 50\text{cm}$  ( $6 \times 10^6 \text{ channels}$ )
- $\sigma_\phi = 16 \mu\text{m}$ ,  $\sigma_z = 580 \mu\text{m}$

## Transition Radiation Tracker (TRT)

- $50 < \text{radii} < 100 \text{ cm}$  ( $4 \times 10^5 \text{ channels}$ )
- $\sigma = 150 \mu\text{m}$  per straw

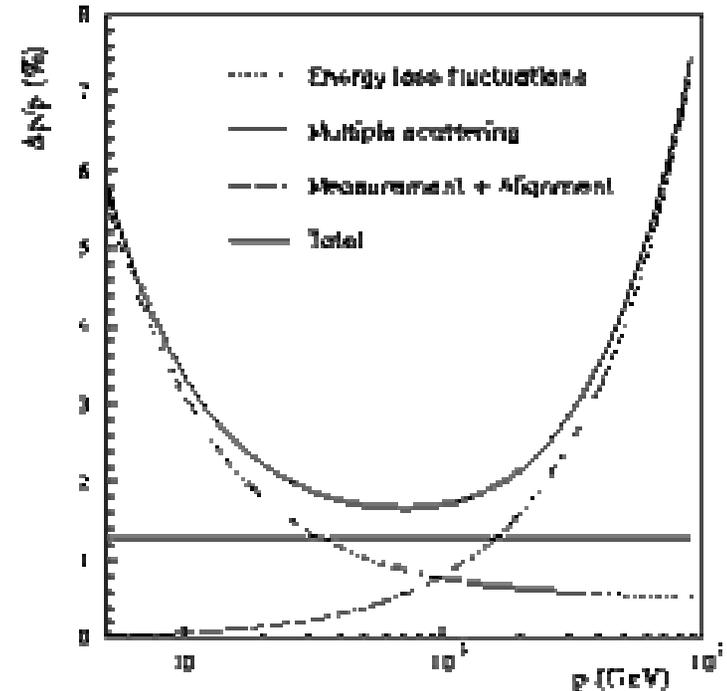
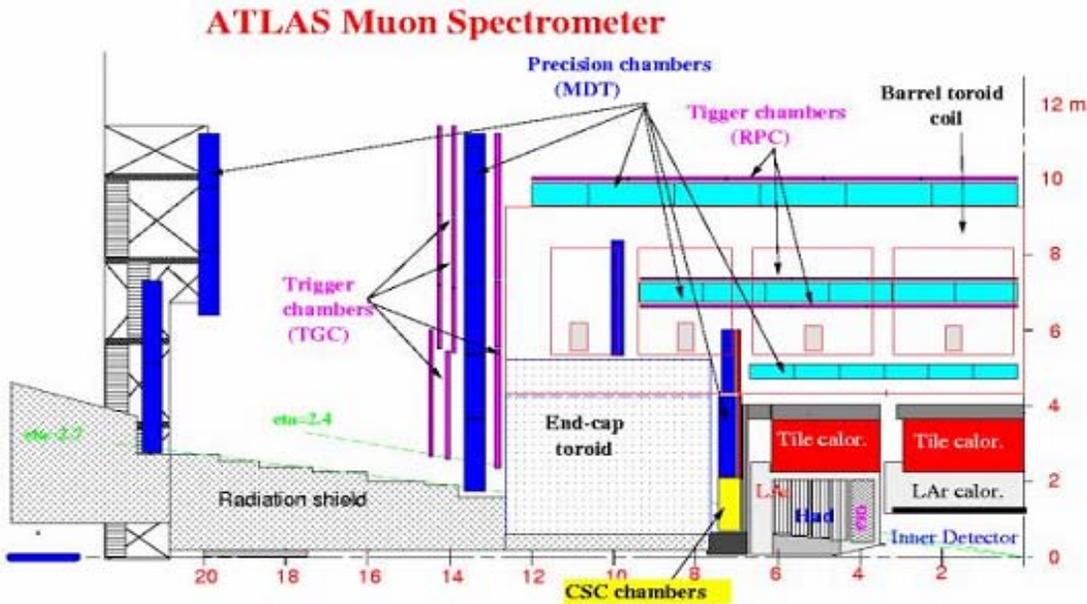
*The silicon detectors provide  $\sim 10$  azimuthal position measurements for  $10 - 20 \mu\text{m}$  resolution.*

*The TRT provides  $\sim 36$  azimuthal position measurements for  $150 \text{ microns}$ .*



# Muon Spectrometer

The momentum of the muons is determined from the curvatures of their tracks in a toroidal magnetic field.



Muon tracks are identified and measured after their passage through  $\sim 2$  m of material.

Track measurement is made with  $\sigma = 60 \mu\text{m}$  intrinsic resolution in three precision measurement (Monitored Drift Tube) stations.

ATLAS is well instrumented for  $B$  Physics: precision vertexing and tracking, good muon id, high resolution calorimetry, and a flexible dedicated  $B$  trigger.

A rich  $B$  Physics program is planned, including CP violation (especially in the  $A_b$  and  $B_s$  systems not accessible to the  $B$  factories), rare decays sensitive to new physics (including  $b \rightarrow s l^+ l^-$  and  $b \rightarrow d l^+ l^-$ ), and baryons and heavy flavor mesons.

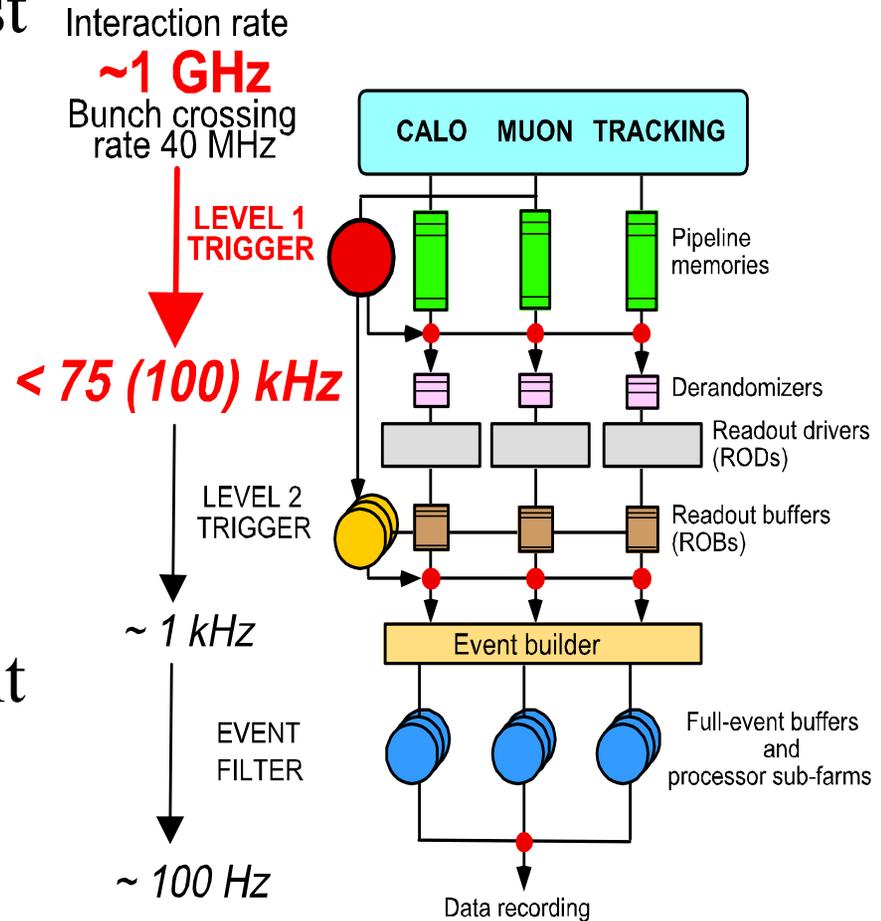
- ATLAS high statistics studies of **heavy flavor hadrons**, including quarkonia, throw light on:
  - bound states
  - the spin dependence of quark confinement
  - the nature of the strong potential
  - Heavy Quark Effective Theory factorization
  - CP violation
  - puzzles in present measurements of hyperon polarization and cross sections.

# Triggering and the B Physics program...

- Muons provide efficient identification and a reliable flavor tag
- **Level 1:** Uses calorimeter + muon trigger chambers, identifies Regions of Interest; outputs 75 kHz. Thresholds are luminosity-dependent. **Paths:**
  - **dimuon** ( $p_T > 6$  GeV [barrel] or  $> 3$  GeV [endcaps])
    - Ex.  $B_d \rightarrow J/\psi K_s^0$ ,  $B_s \rightarrow J/\psi \phi$ ,  $B \rightarrow \mu\mu$ ,  $\Lambda_b \rightarrow \Lambda^0 J/\psi$
  - **MU** ( $p_T > 6$  GeV) + **EM** (cluster  $E_T > 5$  GeV)
    - Ex.  $B_d \rightarrow J/\psi(ee)K_s^0 + b \rightarrow \mu X$ ,  $B_d \rightarrow K^{0*} \gamma$ ,  $B_s \rightarrow \phi \gamma + b \rightarrow \mu X$
  - **MU** ( $p_T > 6$  GeV) + **JET** (cluster  $E_T > 10$  GeV)
    - Ex.  $B_s \rightarrow D_s \phi(KK)$ ,  $B_s \rightarrow D_s(\phi(KK))a_1(\rho^0 \pi^+)$ ,  
 $B^+ \rightarrow K^+ K^+ \pi^-$ ,  $B_d \rightarrow \pi^+ \pi^- (+b \rightarrow \mu X)$

# Trigger, continued

- **Level 2:** uses Regions of Interest and dedicated online algorithms; confirms muons and calorimeter info, refits tracks in Inner Detector; outputs 1 kHz
- **Event Filter:** uses full event buffers + processor subfarms; offline algorithms with alignment and calibration; reconstructs decay vertices; selects exclusive final states via mass and decay length; outputs 200 Hz



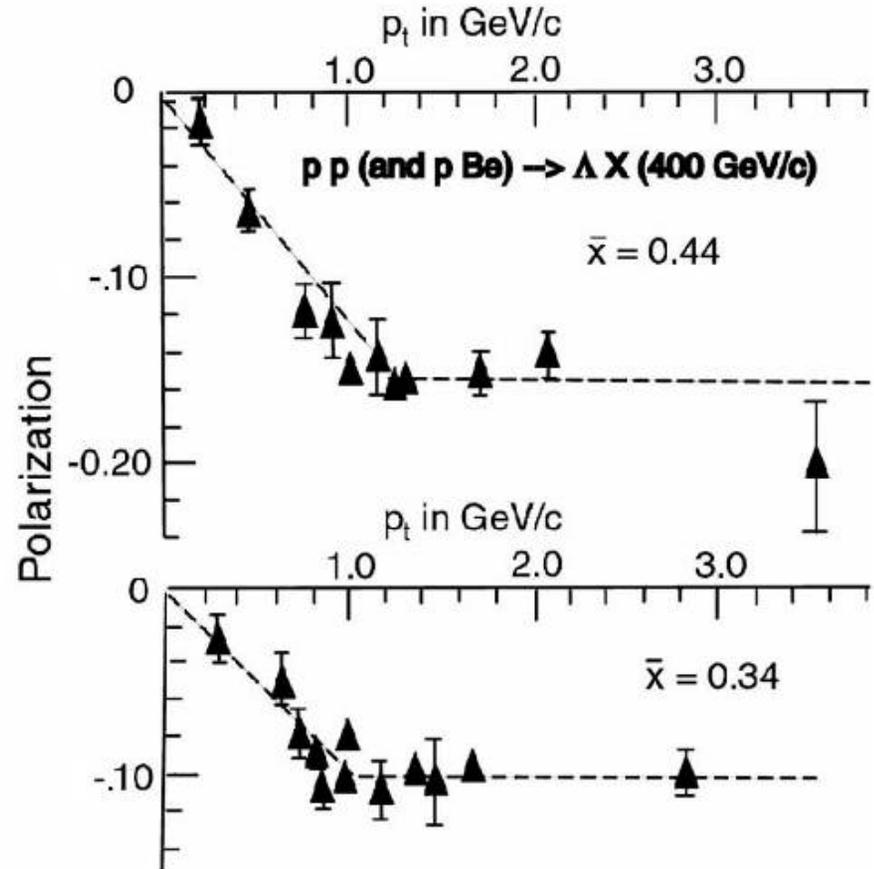
## *Heavy flavor hadron studies underway:*

- Production characteristics, lifetime, and polarization properties of the  $\Lambda_b$
- Production mechanisms, lifetime, and branching ratios for decays in the  $J/\psi$  and  $Y$  systems.
- Resonant and non-resonant decays of  $\chi_b$
- Masses and lifetimes of excited states of the  $B_c$

# $\Lambda_b$ Production, Lifetime, and Polarization

## Motivations:

- What is the role of spin in heavy quark production leading to polarization? What determines the  $\Lambda$  polarization magnitude,  $p_T$  spectrum, and  $x$  dependence? How does polarization depend on quark mass? Are the polarization patterns of  $\Lambda$  and  $\Lambda_b$  the same? Are the  $s$  and  $b$  production mechanisms similar?
- What produced this unusual shape? Expect *small* polarization because the  $\Lambda$ 's are produced inclusively; and most theories predict negligible polarization at high  $p_T$ , high energy.



# $\Lambda_b$ Production, Lifetime, and Polarization

## *Motivations, continued:*

- Longstanding B-baryon lifetime puzzle:  $\tau_{\Lambda_b} / \tau_{B^0}$  data are not consistent with predictions by theory that succeeds for  $\tau_{B^0} / \tau_{B^\pm}$ . Probe this in a new regime with better precision and improved theory.
- Tests perturbative QCD, Heavy Quark Effective Theory, the non-relativistic quark model.
- Probes CP violation in a regime not previously explored: if CP is not conserved, asymmetry parameters  $\alpha_b \neq \alpha_{\bar{b}}$ .

*Goals:* World's first measurements of the asymmetry parameter  $\alpha_b$  and polarization  $P_b$  for  $\Lambda_b$ , to 2%, and the lifetime, to 0.3%.

## *Process:*

- $\Lambda_b \rightarrow \Lambda J/\psi; \Lambda \rightarrow \pi p, J/\psi \rightarrow \mu\mu$

## *Predicting the number of events:*

- Acceptance for  $\Lambda_b$ , from Pythia: 0.157%
- $\sigma(\text{pp} \rightarrow \dots \rightarrow \Lambda_b)$ : 0.00828113 mb
- $BR(\Lambda_b \rightarrow \Lambda J/\psi) \sim 4.7 \times 10^{-4}$
- $BR(\Lambda \rightarrow \pi p) \sim 0.64$
- $BR(J/\psi \rightarrow \mu^+ \mu^-) \sim 0.06$
- Generator-level selection cut efficiency  $\sim 0.05$
- Trigger and reconstruction efficiency  $\sim 0.08$

Anticipated number of  $\Lambda_b$  events in  $30 \text{ fb}^{-1}$ : **18000**

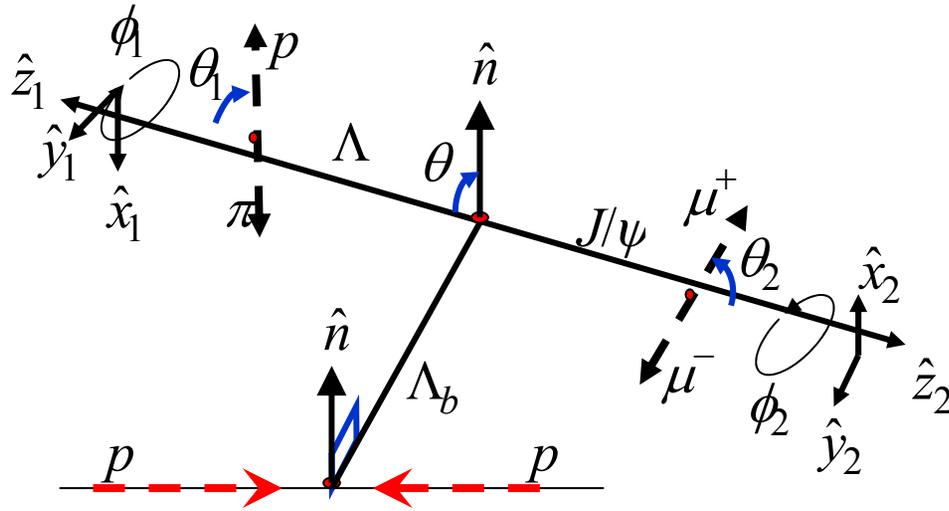
## *Selection Strategy:*

- $p_T > 2.5$  GeV and  $> 4$  GeV for two muons
- $p_T > 0.5$  GeV for  $\pi$  and p
- $\eta < 2.7$  for all tracks

Expected number of background  $pp \rightarrow J/\psi \Lambda X$ :  $3.2 \times 10^6$

A lifetime cut will be used to remove the  $pp \rightarrow J/\psi X$  background.

The **polarization measurement** analyzes the angular distribution of the decay  $\Lambda_b \rightarrow J/\psi(\mu\mu)\Lambda(p\pi)$



The general decay amplitude is given by

$$M = \bar{\Lambda}(p_{\Lambda}) \varepsilon_{\mu}^{*}(p_{J/\psi}) \left[ A_1 \gamma^{\mu} \gamma^5 + A_2 \frac{p_{\Lambda_b}^{\mu}}{m_{\Lambda_b}} \gamma^5 + B_1 \gamma^{\mu} + B_2 \frac{p_{\Lambda_b}^{\mu}}{m_{\Lambda_b}} \right] \Lambda_b(p_{\Lambda_b})$$

where  $A_1$ ,  $A_2$ ,  $B_1$ , and  $B_2$  are model-dependent parameters.

Define:  $a_+ = |a_+|e^{i\alpha_+} \equiv M_{+1/2,0}$        $a_- = |a_-|e^{i\alpha_-} \equiv M_{-1/2,0}$   
 $b_+ = |b_+|e^{i\beta_+} \equiv M_{-1/2,-1}$        $b_- = |b_-|e^{i\beta_-} \equiv M_{+1/2,+1}$

where  $M_{\lambda_1,\lambda_2}$  is the amplitude for decay into  $\Lambda$  with helicity  $\lambda_1$  and  $J/\psi$  with helicity  $\lambda_2$ .

Then for normalized amplitudes,  $\alpha_b = |a_+|^2 - |a_-|^2 + |b_+|^2 - |b_-|^2$

The angular distribution for  $\Lambda_b \rightarrow J/\psi(\mu^+\mu^-)\Lambda(p\pi)$  is

$$w(\vec{A}, \vec{\alpha}, \vec{\theta}) = \sum_{k=0}^{19} f_{1k}(\vec{A}) f_{2k}(\vec{\alpha}) F_k(\vec{\theta})$$

where...

Parameters  $f_{1i}$ ,  
 $f_{2i}$ , and  $F_i$  are:

i	$f_{1i}$	$f_{2i}$	$F_i$
0	$a_+a_+^* + a_-a_-^* + b_+b_+^* + b_-b_-^*$	1	1
1	$a_+a_+^* - a_-a_-^* + b_+b_+^* - b_-b_-^*$	$P_b$	$\cos \theta$
2	$a_+a_+^* - a_-a_-^* - b_+b_+^* + b_-b_-^*$	$\alpha_\Lambda$	$\cos \theta_1$
3	$a_+a_+^* + a_-a_-^* - b_+b_+^* - b_-b_-^*$	$P_b\alpha_\Lambda$	$\cos \theta \cos \theta_1$
4	$-a_+a_+^* - a_-a_-^* + \frac{1}{2}b_+b_+^* + \frac{1}{2}b_-b_-^*$	1	$d_{00}^2(\theta_2)$
5	$-a_+a_+^* + a_-a_-^* + \frac{1}{2}b_+b_+^* - \frac{1}{2}b_-b_-^*$	$P_b$	$d_{00}^2(\theta_2) \cos \theta$
6	$-a_+a_+^* + a_-a_-^* - \frac{1}{2}b_+b_+^* + \frac{1}{2}b_-b_-^*$	$\alpha_\Lambda$	$d_{00}^2(\theta_2) \cos \theta_1$
7	$-a_+a_+^* - a_-a_-^* - \frac{1}{2}b_+b_+^* - \frac{1}{2}b_-b_-^*$	$P_b\alpha_\Lambda$	$d_{00}^2(\theta_2) \cos \theta \cos \theta_1$
8	$-3\text{Re}(a_+a_-^*)$	$P_b\alpha_\Lambda$	$\sin \theta \sin \theta_1 \sin^2 \theta_2 \cos \phi_1$
9	$3\text{Im}(a_+a_-^*)$	$P_b\alpha_\Lambda$	$\sin \theta \sin \theta_1 \sin^2 \theta_2 \sin \phi_1$
10	$-\frac{3}{2}\text{Re}(b_-b_+^*)$	$P_b\alpha_\Lambda$	$\sin \theta \sin \theta_1 \sin^2 \theta_2 \cos(\phi_1 + 2\phi_2)$
11	$\frac{3}{2}\text{Im}(b_-b_+^*)$	$P_b\alpha_\Lambda$	$\sin \theta \sin \theta_1 \sin^2 \theta_2 \sin(\phi_1 + 2\phi_2)$
12	$-\frac{3}{\sqrt{2}}\text{Re}(b_-a_+^* + a_-b_+^*)$	$P_b\alpha_\Lambda$	$\sin \theta \cos \theta_1 \sin \theta_2 \cos \theta_2 \cos \phi_2$
13	$\frac{3}{\sqrt{2}}\text{Im}(b_-a_+^* + a_-b_+^*)$	$P_b\alpha_\Lambda$	$\sin \theta \cos \theta_1 \sin \theta_2 \cos \theta_2 \sin \phi_2$
14	$-\frac{3}{\sqrt{2}}\text{Re}(b_-a_-^* + a_+b_+^*)$	$P_b\alpha_\Lambda$	$\cos \theta \sin \theta_1 \sin \theta_2 \cos \theta_2 \cos(\phi_1 + \phi_2)$
15	$\frac{3}{\sqrt{2}}\text{Im}(b_-a_-^* + a_+b_+^*)$	$P_b\alpha_\Lambda$	$\cos \theta \sin \theta_1 \sin \theta_2 \cos \theta_2 \sin(\phi_1 + \phi_2)$
16	$\frac{3}{\sqrt{2}}\text{Re}(a_-b_+^* - b_-a_+^*)$	$P_b$	$\sin \theta \sin \theta_2 \cos \theta_2 \cos \phi_2$
17	$-\frac{3}{\sqrt{2}}\text{Im}(a_-b_+^* - b_-a_+^*)$	$P_b$	$\sin \theta \sin \theta_2 \cos \theta_2 \sin \phi_2$
18	$\frac{3}{\sqrt{2}}\text{Re}(b_-a_-^* - a_+b_+^*)$	$\alpha_\Lambda$	$\sin \theta_1 \sin \theta_2 \cos \theta_2 \cos(\phi_1 + \phi_2)$
19	$-\frac{3}{\sqrt{2}}\text{Im}(b_-a_-^* - a_+b_+^*)$	$\alpha_\Lambda$	$\sin \theta_1 \sin \theta_2 \cos \theta_2 \sin(\phi_1 + \phi_2)$

The angular distribution depends on 9 unknown parameters:  $P_b$  and 4 amplitudes and 4 phases of the  $a_+$ ,  $a_-$ ,  $b_+$ , and  $b_-$ .

After normalization  $|a_+|^2 + |a_-|^2 + |b_+|^2 + |b_-|^2 = 1$  and global phase constraint, number of independent unknowns=7.

Extract these from the measured decay angles by a five-dimensional likelihood fit:

$$L = -2 \sum_{j=1}^N \log(w_{obs}(\vec{\theta}', \vec{A}, P_b))$$

where  $N =$  number of events,

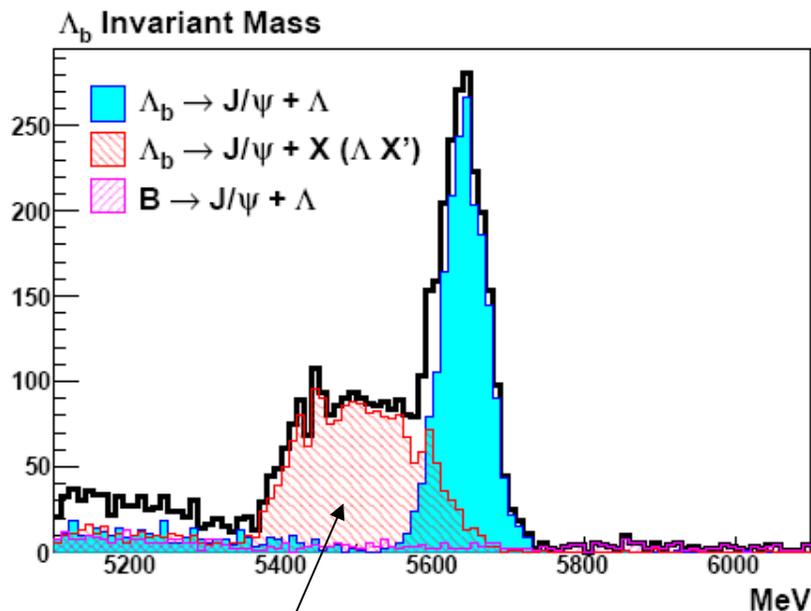
$$w_{obs}(\vec{\theta}', \vec{A}, P_b) = \frac{\int w(\vec{\theta}, \vec{A}, P_b) T(\vec{\theta}, \vec{\theta}') d\vec{\theta}}{\iint w(\vec{\theta}, \vec{A}, P_b) T(\vec{\theta}, \vec{\theta}') d\vec{\theta} d\vec{\theta}'}$$

$$T(\vec{\theta}, \vec{\theta}') = \varepsilon(\vec{\theta}) R(\vec{\theta}, \vec{\theta}')$$

for  $\varepsilon =$  acceptance and  $R =$  resolution.

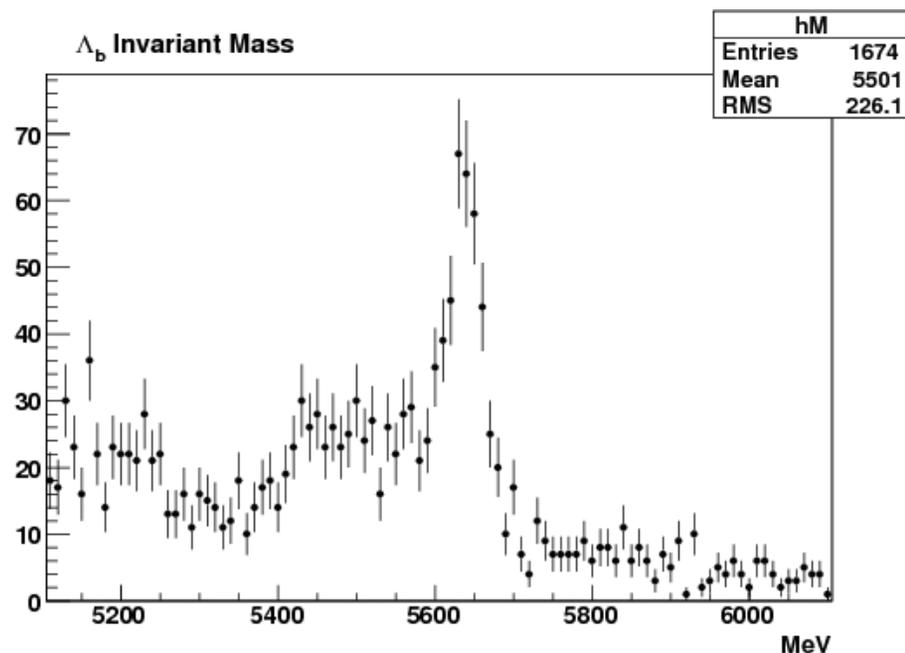
The resolution on the angles measured is  $\sim 10$  mrad. Detector acceptance corrections come from a phase space  $\Lambda_b$  monte carlo sample.

# Expected signal and background shapes, using Pythia with EVTGEN:



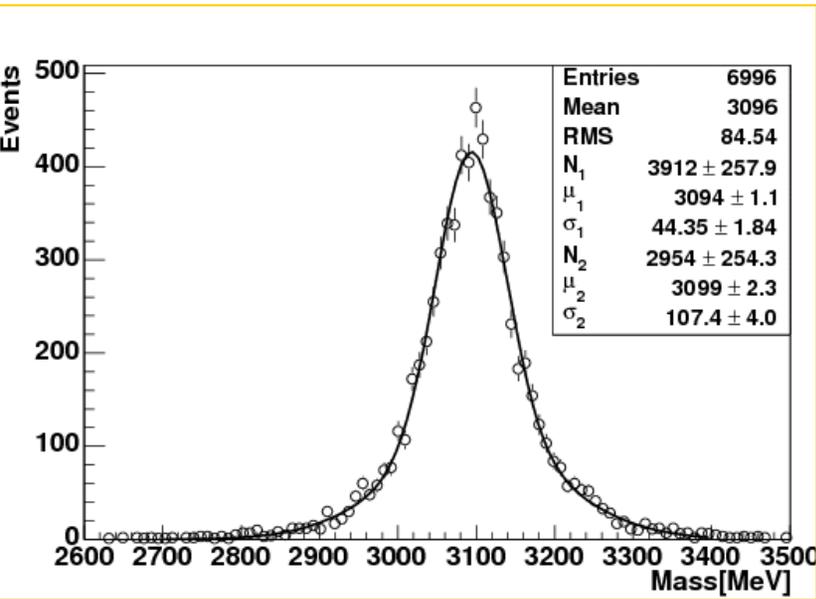
Generated and smeared

Note this enhancement reflects all  $\Lambda_b \rightarrow J/\psi X$  BR's set to same BR as  $\Lambda_b \rightarrow J/\psi \Lambda$  --- unlikely situation in the real data.

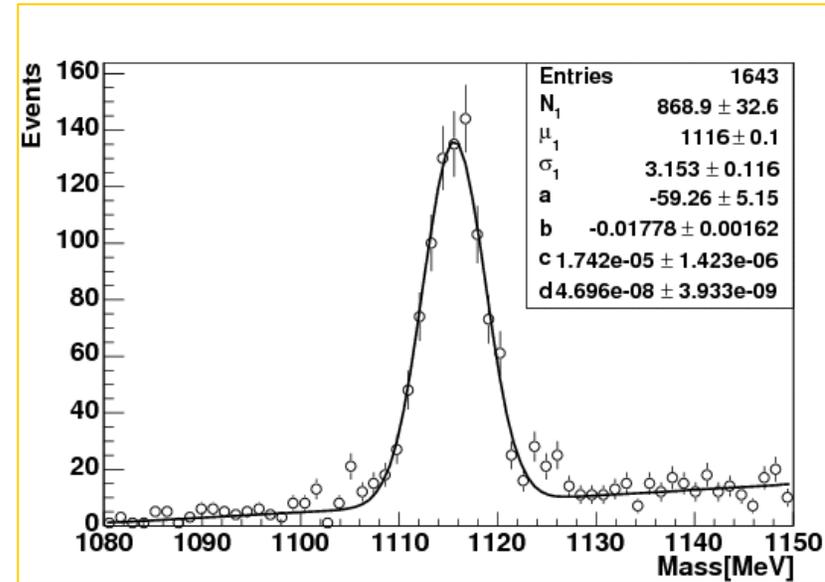


Reconstructed, excluding trigger effects

# Mass distributions

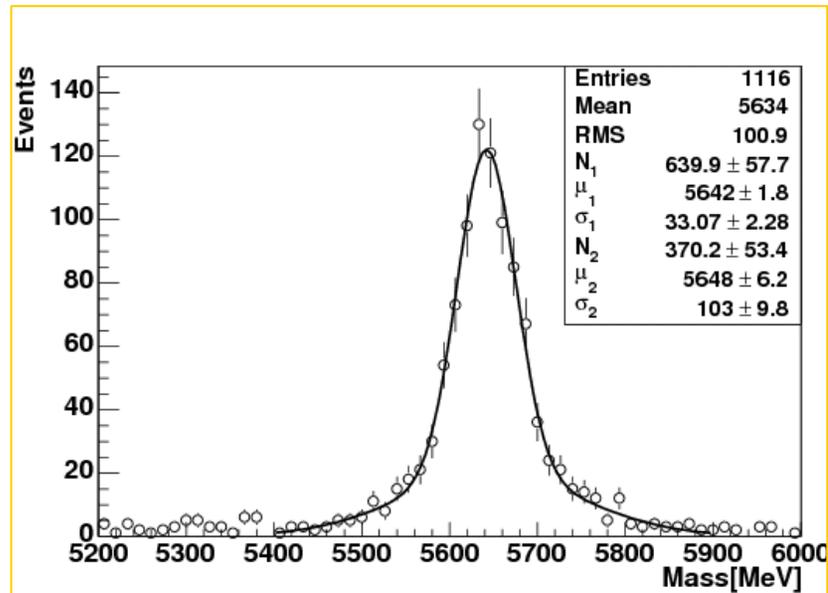


$J/\psi \rightarrow \mu\mu$ , with  $\mu\mu$  vertexing



$\Lambda \rightarrow \pi p$ , after  $J/\psi \Lambda$  vertex required

$\Lambda_b \rightarrow J/\psi \Lambda$ , with vertexing and  $J/\psi$  and  $\psi$  mass constraints



## *A word about heavy quarkonium production models...*

One of the goals of the next two studies is a comparison of ATLAS data to predictions by the quarkonium production models---esp. the **Color Singlet Model** (*assumes each quarkonium state can be produced only by a  $q\bar{q}$  pair in the same color and  $J$  state as the quarkonium*) and Non-Relativistic QCD with the **Color Octet Mechanism** (*treats quarkonium as a non-relativistic system; thus  $q\bar{q}$  pairs produced with one set of quantum numbers can evolve into a quarkonium state with different quantum numbers by emitting low energy gluons.*)

# Heavy quarkonia production and decay

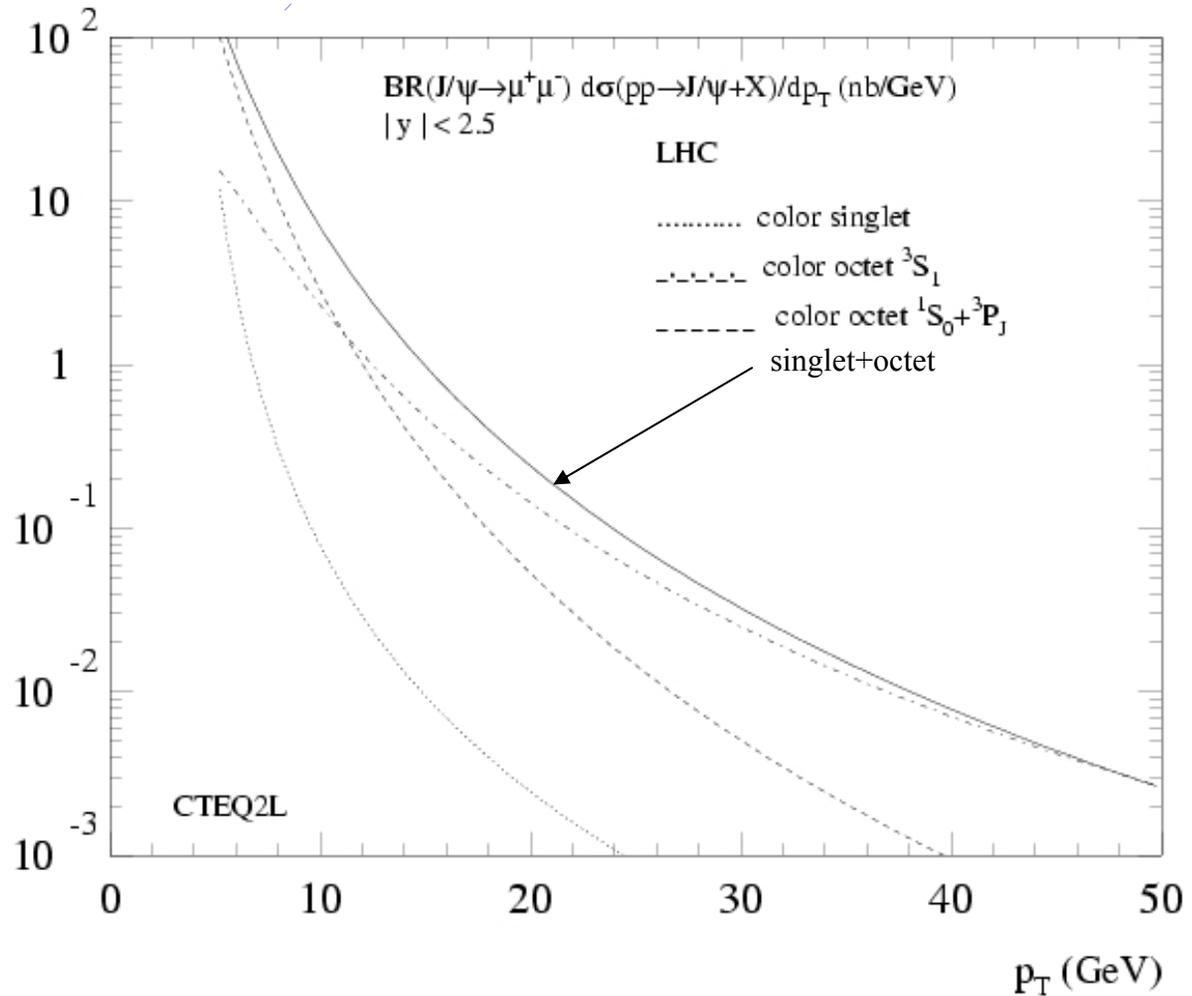
## *Motivation:*

- The production mechanism of quarkonia is not fully understood. Several models have been proposed. NRQCD with the Color Octet Mechanism, with free normalization, agrees with CDF data on  $J/\psi$   $p_T$  cross section up to accessible Tevatron energies. Its polarization vs.  $p_T$  predictions have not been confirmed. Additional data, especially at higher  $p_T$  and for other onia farther from the QCD scale, are needed.

## *Goals:*

- Distributions in  $p_T$ ,  $\eta$ , and polarization for  $J/\psi$ ; measurement of the ratio of  $Y$  and  $J/\psi$  production cross sections; assessment of hadronic activity associated with quarkonium production.

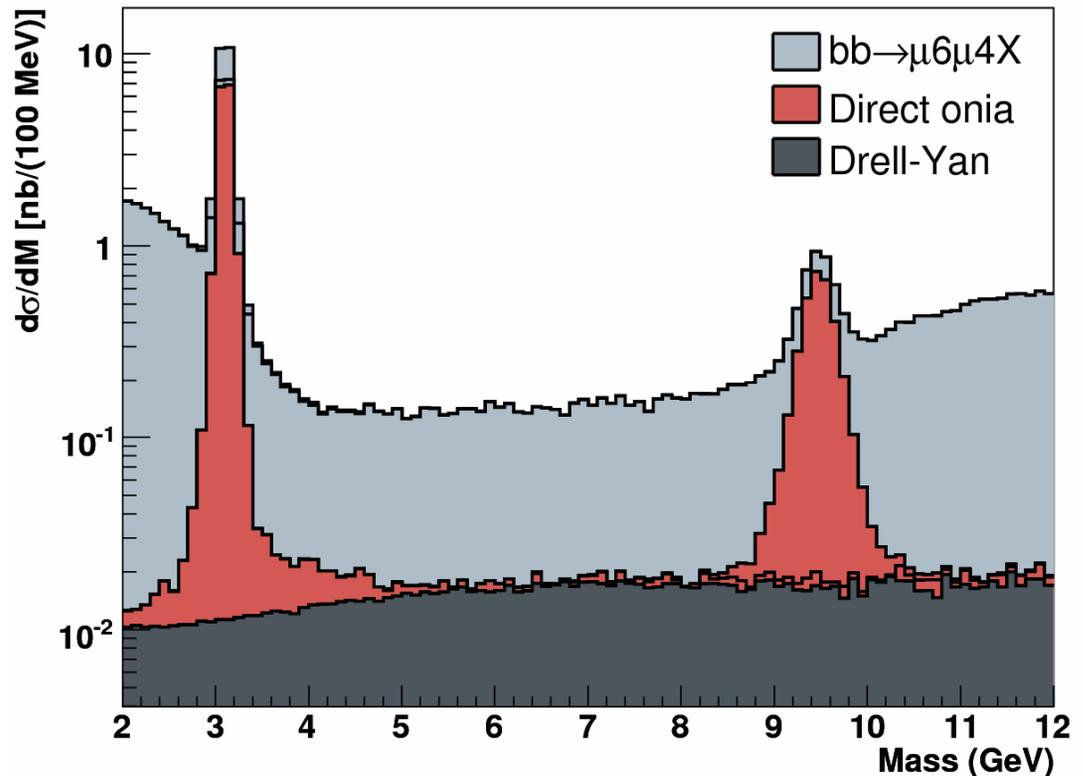
The LHC will produce heavy quarkonia with high  $p_T$  at high rate:



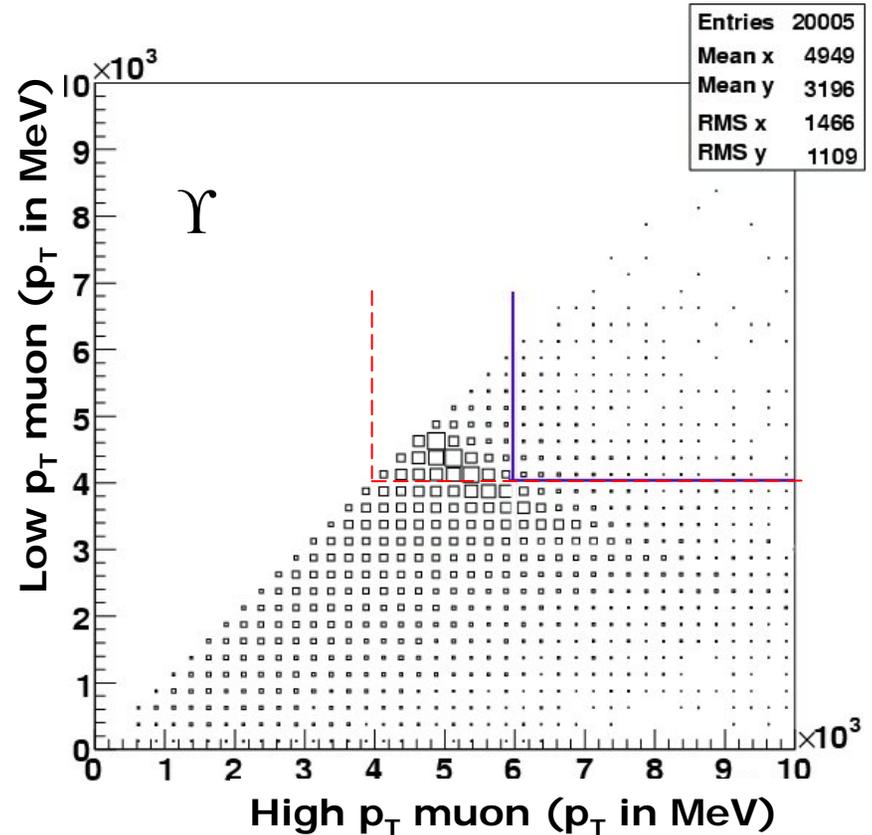
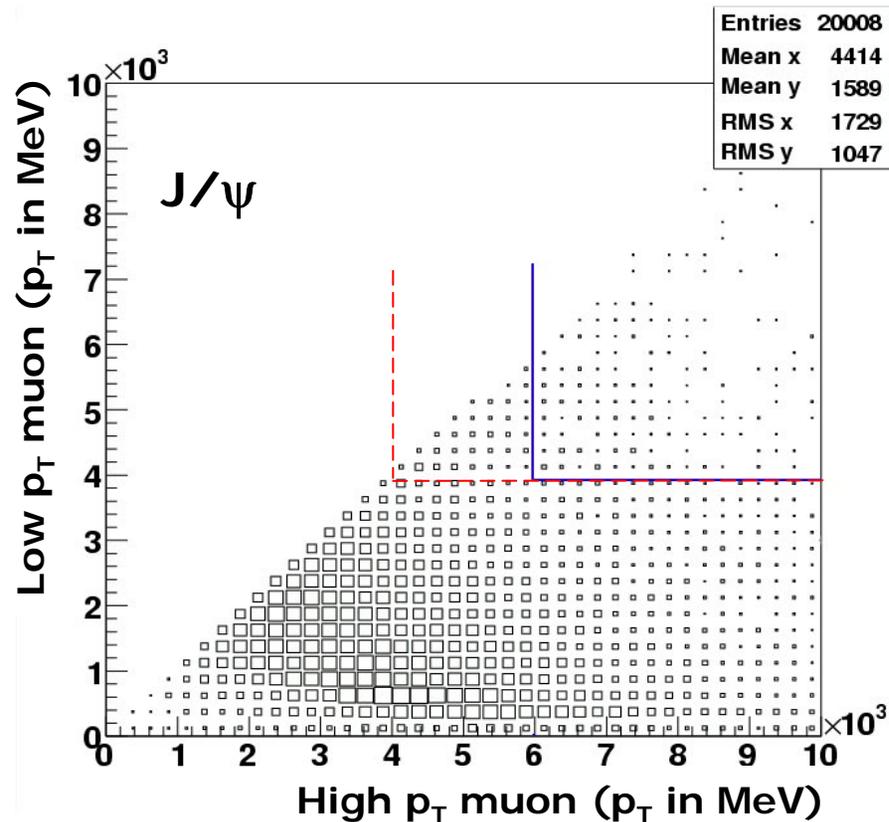
*Processes and #events expected in  $10^6$  seconds  
@  $L=10^{31}/\text{cm}^2\text{s}$ :*

- $J/\psi \rightarrow \mu(p_T \geq 6 \text{ GeV}) \mu(p_T \geq 4 \text{ GeV})$ : 175k events
- $Y \rightarrow \mu(p_T \geq 6 \text{ GeV}) \mu(p_T \geq 4 \text{ GeV})$ : 36k events

*Generator-level study of  
the backgrounds from  
 $bb \rightarrow \mu\mu X$  and Drell-Yan  
production, including  
trigger and detection  
efficiency:*



**Analysis strategy #1:** Examine the option to lower dimuon  $p_T$  trigger threshold from 6 GeV+4 GeV to 4 GeV+4 GeV

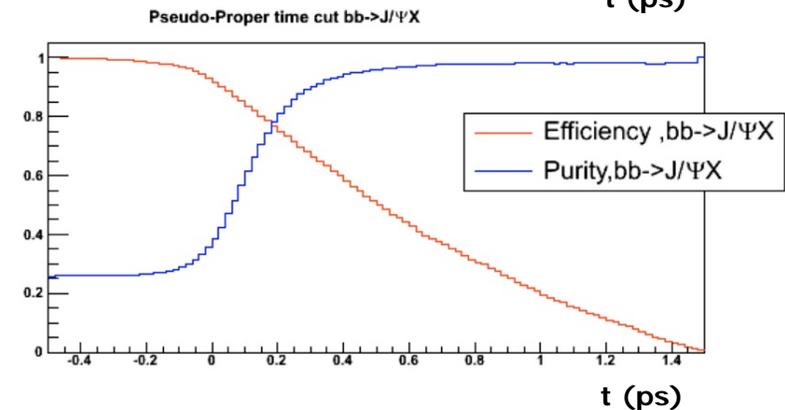
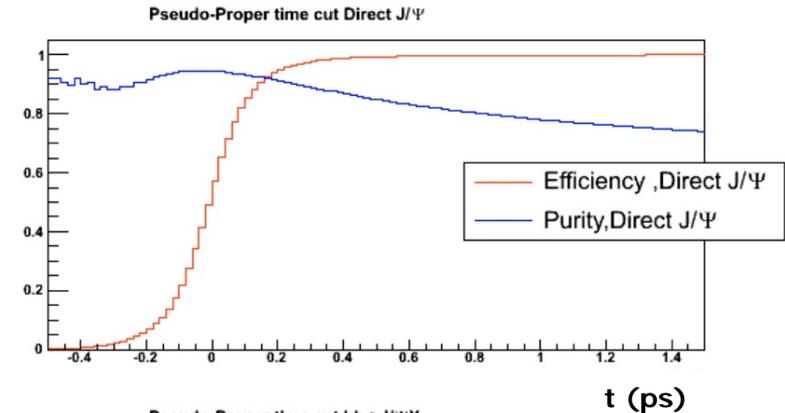
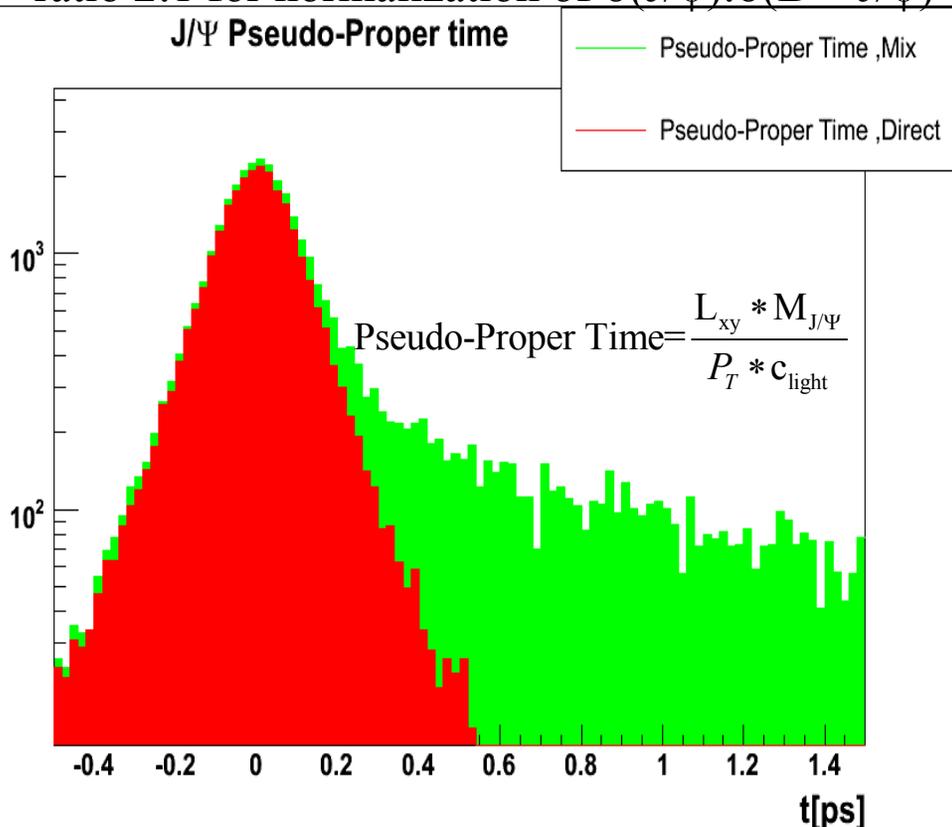


*This also increases the contribution of color singlet production, which dominates for  $p_T < 10$  GeV.*

## Analysis strategy #2:

Cut on pseudo-proper time to separate direct from indirect  $J/\psi$ 's. As expected, the resolution worsens with increasing  $p_T$  and  $\eta$ .

“Mix” sample is direct + indirect  $J/\psi$  combined, with ratio 2:1 for normalization of  $\sigma(J/\psi):\sigma(B \rightarrow J/\psi)$



## Analysis strategy #3: Measure the quarkonium spin

For the angle  $\theta$  between the positive muon in the quarkonium reference frame and the quarkonium direction in the lab frame,

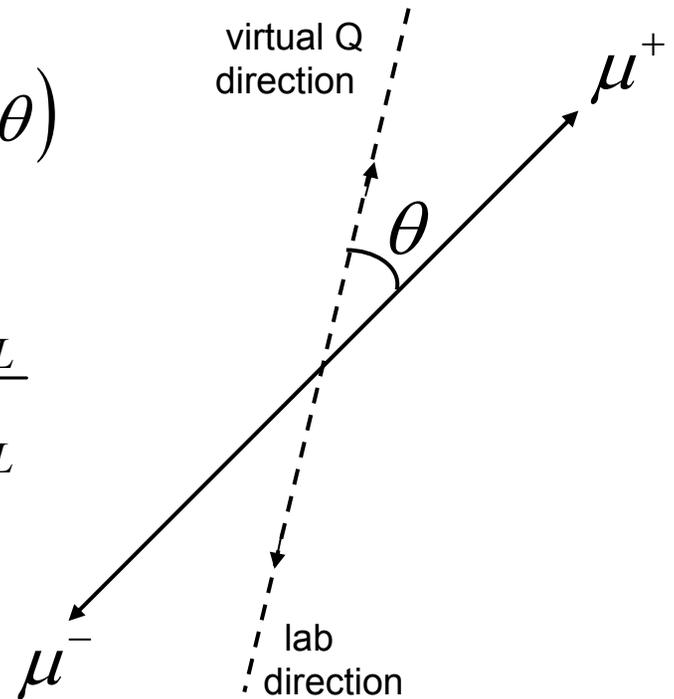
$$\frac{d\Gamma}{d\cos\theta} \propto (1 + \alpha \cos^2 \theta)$$

$$\text{where } \alpha \equiv \frac{\sigma_T + 2\sigma_L}{\sigma_T - 2\sigma_L}$$

$\alpha = 0$ : unpolarized meson

$\alpha = +1$ : transverse polarization

$\alpha = -1$ : longitudinal polarization

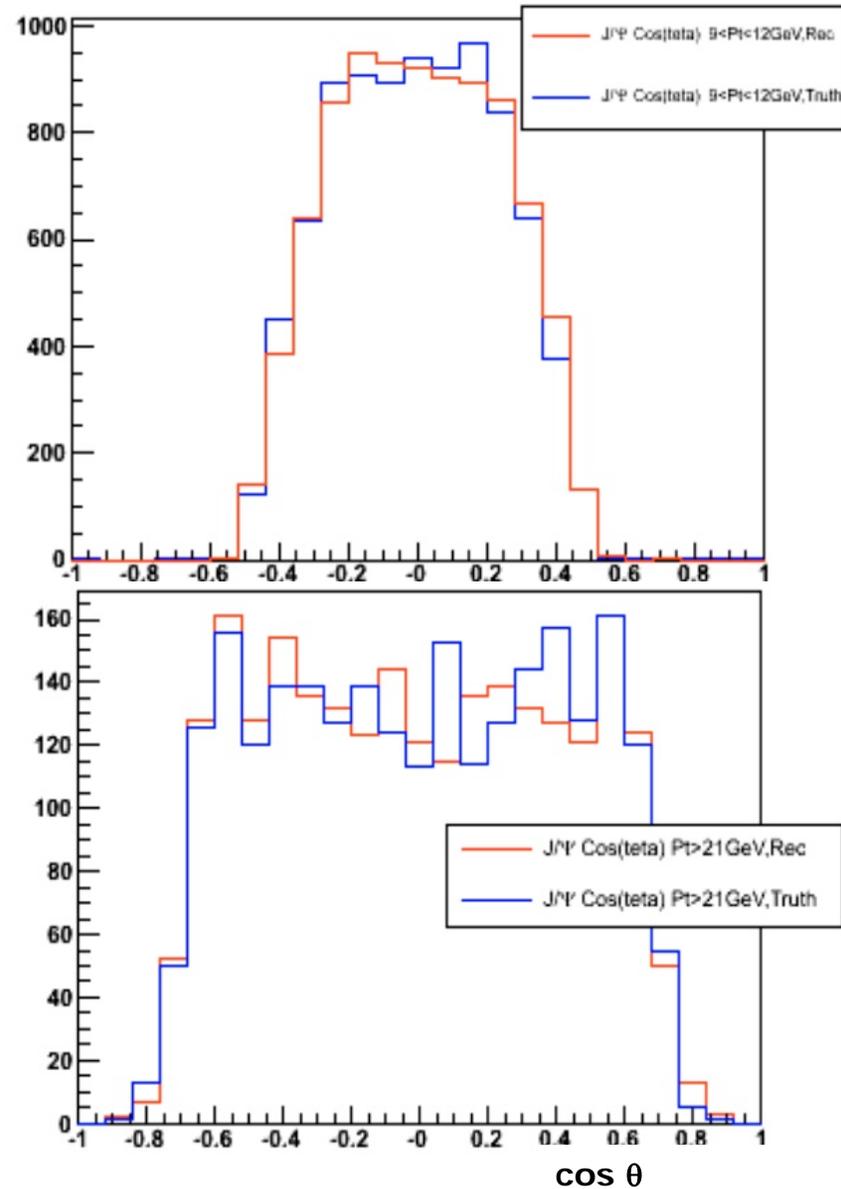


*Octet production predicts transverse polarization at large  $p_T$ .*

*Preparation for this polarization study...*

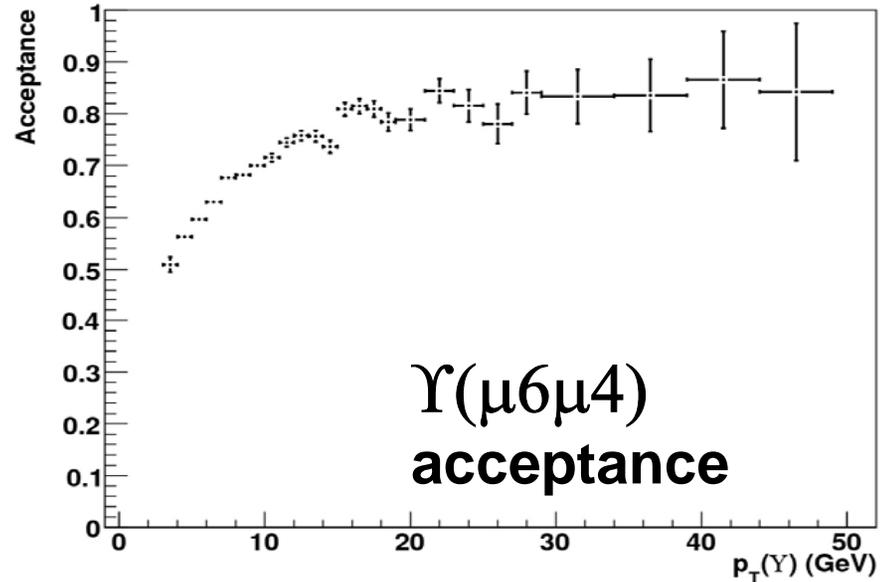
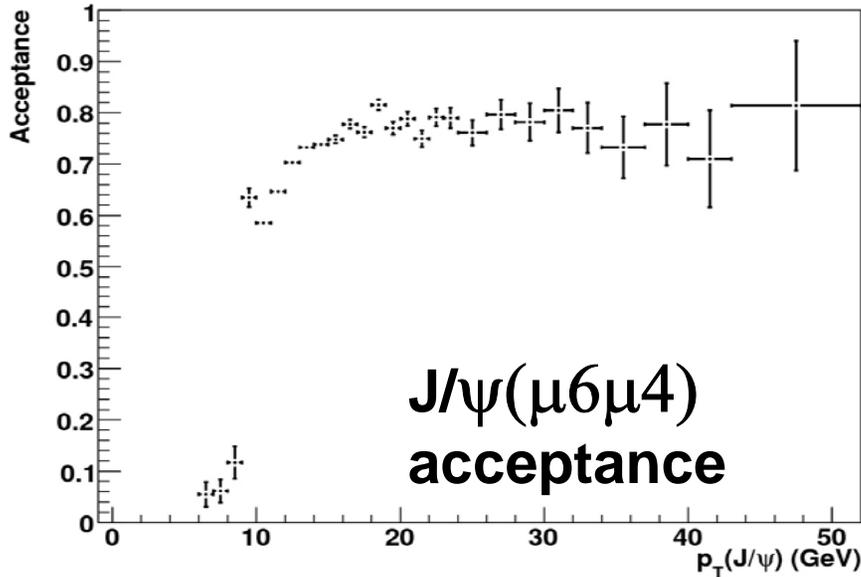
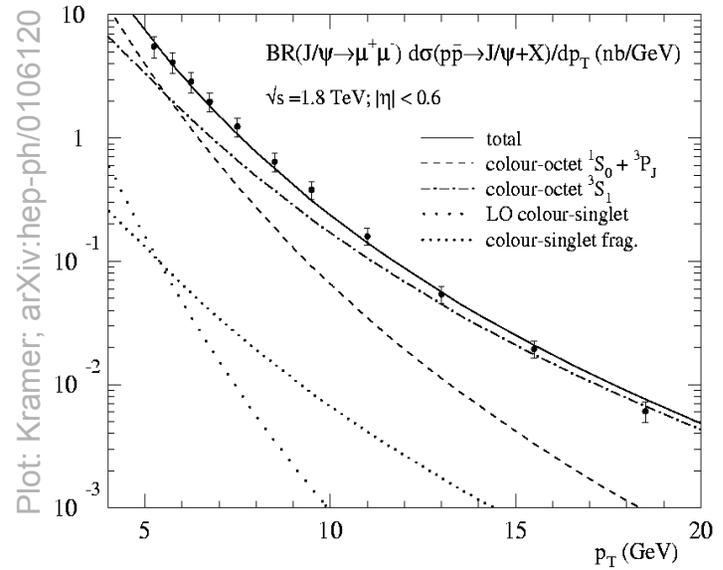
Reconstructed versus truth distributions for simulated  $J/\psi$  events:

- Acceptance in  $\cos \theta$  is limited by the  $p_T$  of the second muon.
- A single muon trigger would be useful if the rate could be reduced by other means.



The present NRQCD formalism models Tevatron data up to  $p_T=20$  GeV well  $\rightarrow$

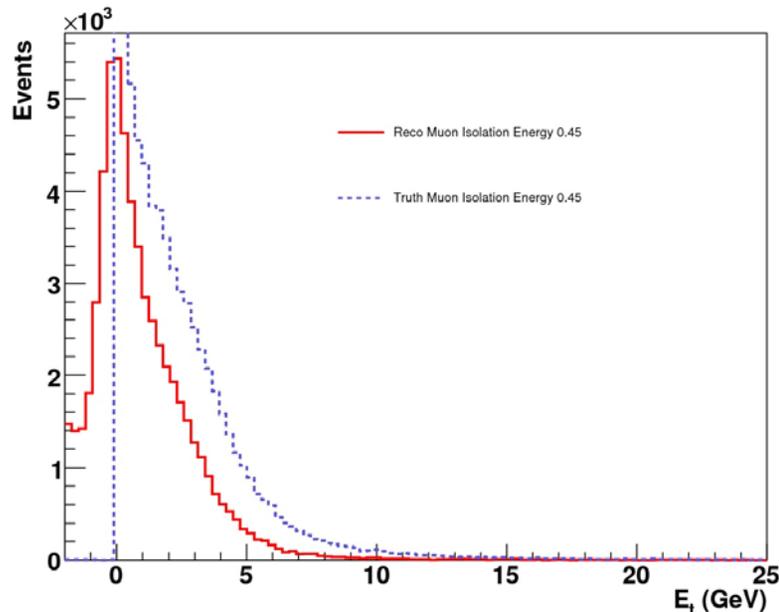
Will this success continue at higher  $p_T$ ? Exploit the high LHC acceptance for onia  $p_T$  up to 50 GeV  $\downarrow$



## Analysis strategy #4: the jet analysis

*The Color Singlet and Color Octet Models may predict different levels of hadronic activity associated with onium production. Prepare to study this with isolation cones about the onium direction.*

First step: study using the isolation cone about the muon direction:



Muons from onium must be subtracted.

But 29% of the  $J/\psi$  cross section comes from  $\chi_c$  decays, which have an associated energetic  $\gamma$ . These photons must be modelled.

For  $J/\psi$  the photon and muon directions are almost collinear.

# $\chi_b$ Resonant and Non-resonant Decay

## *Motivations:*

- A single measurement of  $\text{Br}(\chi_{c0} \rightarrow \phi\phi)$  exists and is substantially larger than prediction. Is this due to internal motion of the quarks in the hard part of the amplitude? Probe this hypothesis for analogous channel  $\chi_{bJ} \rightarrow \psi\psi X$  at LHC.
- The color singlet model for  $\psi$  production via  $gg \rightarrow \psi + g$  is inconsistent with Tevatron data. Color Octet Model parameters can take input from LHC data.

## *Search for process:*

- $\chi_{b0,2} \rightarrow \psi\psi \rightarrow \mu^+ \mu^- \mu^+ \mu^-$

*Predicting the statistics at the LHC, and the expected enhancement relative to the Tevatron:*

For  $L=10^{33}/\text{cm}^2\text{s}$ :

$$\sigma(pp \rightarrow \chi_{bi} + X) = \left\{ \begin{array}{l} 250 \text{ nb at Tevatron} \\ 1.5 \text{ } \mu\text{b at LHC} \end{array} \right\} \text{ for } i=0$$
$$\left\{ \begin{array}{l} 320 \text{ nb at Tevatron} \\ 2.0 \text{ } \mu\text{b at LHC} \end{array} \right\} \text{ for } i=2$$

$$BR(\chi_{bi} \rightarrow \psi\psi) = \left\{ \begin{array}{l} 2.2 \times 10^{-4} \\ 5 \times 10^{-4} \end{array} \right. \begin{array}{l} \text{for } i=0 \\ \text{for } i=2 \end{array}$$

$$BR(J / \psi \rightarrow \mu^+ \mu^-) = 5.93\%$$

## *Selection strategy:*

- for muon  $p_T^{\text{Level 1}} > 6 \text{ GeV}$ ,
- muon  $p_T^{\text{Level 2}} > 4 \text{ GeV}$ ,
- muon  $\eta^{\text{Level 1 and 2}} > 2.5$
- for Pythia subprocesses for  $\chi_b$  production:
  - $gg \rightarrow bb\sim[3P0(1)]+g$
  - $gg \rightarrow bb\sim[3P0(1)]+q$ ,

*Efficiency = 1.63%.*

Number of expected  $\chi_{b0}$  events per year: **200.**

Background (4  $\mu$ 's forming 2  $J/\psi$ 's): **almost 0.**

# B<sub>c</sub> Production and Decay

*Motivation:*

Precision reconstruction of B<sub>c</sub> ground and excited states can be used to constrain strong potential models.

*Analysis strategy:*

• Select hadronic B<sub>c</sub> decays. The expected mass difference between 2S and 1S is  $\sim 600$  MeV. Possible channels:

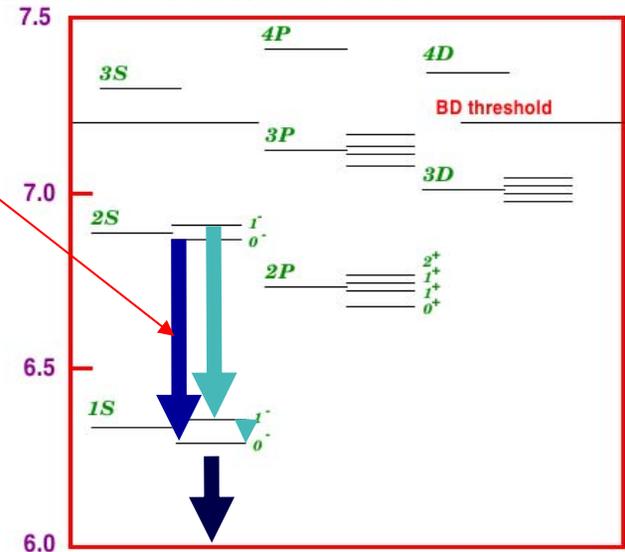
$$B_c^*(2^1S_0) \rightarrow B_c(1^1S_0) \pi^+ \pi^-$$

$$B_c(1^1S_0) \rightarrow J/\psi \pi^+$$

$$B_c^*(2^1S_1) \rightarrow B_c^*(1^1S_1) \pi^+ \pi^-$$

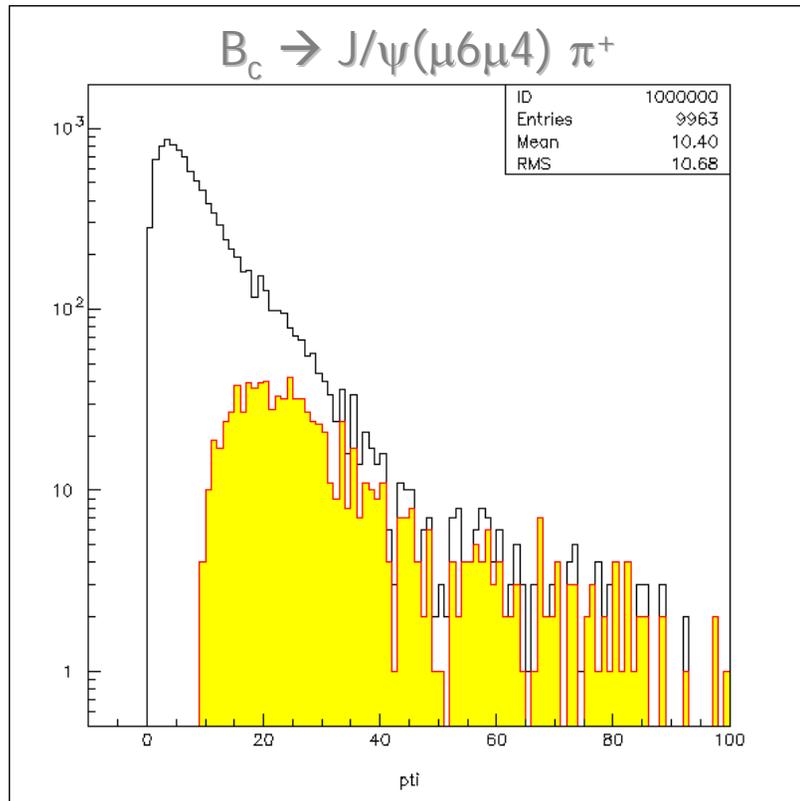
$$B_c^*(1^1S_1) \rightarrow B_c(1^1S_0) + \gamma$$

$$B_c(1^1S_0) \rightarrow J/\psi \pi^+$$



The  $p_T$  of  $B_c$  using standalone generator BCVEGPY2.1 (white)

and  $p_T$  of  $B_c$  whose muons pass the 6 GeV + 4 GeV requirement (yellow):



Needed:

- Efficient tracking at very low  $p_T$
- Excited state generation and decay in Pythia

*Expected number of  $B_c$  events in  $20 \text{ fb}^{-1}$ :*

$$\#b\bar{b} \text{ pairs} \times P(b \rightarrow B_c) \times \text{BR}(B_c \rightarrow J/\psi\pi) \times \text{BR}(J/\psi \rightarrow \mu\mu) \times \varepsilon(\mu\mu\text{-trig}) \times \varepsilon(\text{cuts}):$$
$$1.4 \times 10^{11} \times 0.0015 \times 0.002 \times 0.0593 \times 0.61 \times 0.64 =$$

**10000 events.**

# Conclusions

- With  $1 \text{ fb}^{-1}$ , the  $\Lambda_b$  studies will yield:
  - the  $\Lambda_b$  lifetime, known better than the present world average
  - first results on the  $\Lambda_b$  polarization.
- The large predicted onia cross sections at the LHC will permit definitive studies of:
  - quarkonium spin alignment as a test of the Color Octet Mechanism.  $100 \text{ pb}^{-1}$  are needed for a competitive polarization measurement.
  - jet activity associated with onium decay: the Color Singlet and Color Octet Mechanisms predict different associated jet activity levels.

## Conclusions, *continued*

- with  $10 \text{ pb}^{-1}$ , ATLAS will measure ratios of onia cross sections to constrain NRQCD octet matrix elements. Subsequent statistics will fix the matrix elements.
- the LHC will be the first opportunity for significant statistics on excited states in the  $B_c$  family. These measurements can constrain models of the strong potential and cast light on interaction between the EW and strong forces.
- $\chi_b$  decays may point the way to solving longstanding puzzles in the  $gg \rightarrow \psi + g$  and  $e^+e^- \rightarrow \psi\eta_c$  cross sections.