

Physics 301

Homework due 14 September 2022

- 1) Find an expression for the density of states of a relativistic, massless, classical particle confined to the range $0 < r < r_0$ and $0 < p < p_0$.
- 2) Stowe problem 3-8.
- 3) Consider a system of 100 air molecules in an otherwise empty room. Calculate the probability that 49 are in the front 60% of the room and 55 are in the top half.
- 4) Give an example of a process that involves continuous change but is
 - (a) quasi-static
 - (b) not quasi-static

Support both answers numerically. Choose examples that are different from those in the Stowe book.

- 5) How would the H-Theorem be changed for a system in which the transition rate W_{rs} between states r and s were not equal to the rate for the reverse process?

Answers to homework due 14 Sept. 2022

① General expression for the

quantum state in a finite volume $= \int_V \int_P \frac{1}{h^3} dV_r dV_p$

$$= \int_P \frac{4\pi V_r}{h^3} p^2 dp$$

Let $E = pc$, so

$$p^2 = \frac{E^2}{c^2} \text{ and}$$

$$dp = \frac{1}{c} dE$$



$$= \int_0^{c p_0} \frac{4\pi V_r}{c^3 h^3} E^2 dE$$

$$g(E)$$

$$(2) a) P_N(n) = \frac{N!}{n!(N-n)!} p^n q^{N-n}$$

Here $p = \text{prob of 4 dots up} = \frac{1}{6}$

$q = \text{prob of not 4 dots up} = \frac{5}{6}$

$N = \# \text{ elements (dice)} = 8$

$n = \# \text{ of elements satisfying criterion "p"} = 5$

Plug in

$$P_8(5) = \frac{8!}{5!3!} \left(\frac{1}{6}\right)^5 \left(\frac{5}{6}\right)^3 = \boxed{0.00417}$$

b) The number of configurations is given by $\frac{N!}{n!(N-n)!}$

$$\text{Here the number is } \frac{8!}{5!3!} = \boxed{56}$$

c) Again use $P_N(n) = \frac{N!}{n!(N-n)!} p^n q^{N-n}$

= but now add the cases for $n=5, 6, 7, 8$

$$P_{\text{TOT}} = \sum_{n=5}^8 \frac{8!}{n!(8-n)!} \left(\frac{1}{6}\right)^n \left(\frac{5}{6}\right)^{8-n} = 0.00417$$

$$+ 0.00042$$

$$+ 0.00002$$

$$0.$$

$$\boxed{0.00461}$$

3. Using Stove equation 3.14,

$$P_{100}(49, 55) = P_{100}(49) \cdot P_{100}'(55).$$

Each term is separately calculated with equation 3.12,

$$P_N(n) = \frac{N!}{n!(N-n)!} p^n q^{N-n}. \text{ Plugging in,}$$

$$P_{100}(49, 55) = \frac{100!}{49!51!} (0.6)^{49} (0.4)^{51} \times \frac{100!}{55!45!} (0.5)^{55} (0.5)^{45}$$

To find the factorials, use Stirling's approximation,

$$\ln(m!) \approx m \ln m - m + \frac{1}{2} \ln(2\pi m)$$

$$\text{So } \ln(100!) = 100 \ln 100 - 100 + \frac{1}{2} \ln(2\pi \cdot 100) = 363.6$$

$$\text{So } 100! \approx e^{364}$$

$$\text{Similarly, } 49! \approx e^{144}$$

$$51! \approx e^{152}$$

$$55! \approx e^{168}$$

$$45! \approx e^{129}$$

$$\text{Also } (0.6)^{49} = 1.35 \times 10^{-11}$$

$$(0.4)^{51} = 5.07 \times 10^{-21}$$

$$(0.5)^{100} = 7.89 \times 10^{-31}$$

$$\text{So } P_{100}(49, 55) = e^{(364 + 364 - 144 - 152 - 168 - 129)} \times 1.35 \times 5.07 \times 7.89 \times 10^{-63}$$

$$= \boxed{0.0023}$$

4) A process is quasi-static if it occurs over a time period that is long compared to the time it takes the system to return to equilibrium after being perturbed by the process.

Some orders-of-magnitude relaxation time τ are:

• for quantum mechanical transitions, $\Delta E \cdot \Delta t \approx \hbar$, so

$$\tau \approx \frac{\hbar}{\Delta E} = \frac{10^{-15} \text{ eV} \cdot \text{sec}}{10 \text{ eV}} \approx 10^{-16} \text{ sec}$$

(5)

For air in a room: order of magnitude estimates:

10^{28} molecules

Room volume = 60 m^3

mean distance between molecules: $2 \times 10^{-9} \text{ m}$

Room temperature = 300 K

molecule mass $\approx 3 \times 10^{-26} \text{ kg}$

$$\text{Typical particle velocity} = v = \sqrt{\frac{3kT}{m}} = \sqrt{\frac{3 \cdot (1.38 \times 10^{-23} \text{ J/K}) \cdot 300}{3 \times 10^{-26} \text{ kg}}}$$

$$= 643 \text{ m/s}$$

Suppose that a perturbation occurs in 1 corner of the room; we want to estimate how long it takes for the effect to propagate to the opposite corner.

Typical time for transport of the perturbation:

$$\frac{1}{\sqrt{2}} \frac{\sqrt{60 \text{ m}^3}}{643 \text{ m/s}} = 4 \times 10^{-3} \text{ sec.}$$

So if the air in a room is compressed during a time significantly shorter than 10^{-3} sec , the compression is not quasi-static; if the compression period is significantly longer than 10^{-3} sec , the process is quasi-static.

(6)

5.) We redo the H Theorem for the case where $W_{sr} \neq W_{rs}$. We follow the format used in the handout given to the class.

Consider P_r , the probability that a system is in state r .

If any of the other states of the system can be labelled s ,

then the change per unit time in the probability of being in state r is

$$\frac{dP_r}{dt} = \sum_s P_s W_{sr} - \sum_s P_r W_{rs}, \text{ where}$$

W_{rs} = transition rate out of r and

W_{sr} = " " into r

Define $H \equiv \sum_r P_r \ln P_r$. Then

$$\frac{dH}{dt} = \sum_r \left(\frac{dP_r}{dt} \ln P_r + P_r \frac{d \ln P_r}{dt} \right) = \sum_r \frac{dP_r}{dt} (\ln P_r + 1)$$

Eq 1" $\frac{dH}{dt} \downarrow = \sum_r \sum_s (P_s W_{sr} - P_r W_{rs}) (\ln P_r + 1)$

Rewrite Eq 1 with the summation indices exchanged:

Eq 2" $\frac{dH}{dt} = \sum_s \sum_r (P_r W_{rs} - P_s W_{sr}) (\ln P_s + 1)$

Add Eq 1 + Eq 2:

$$\frac{dH}{dt} = \sum_r \sum_s (P_s W_{sr} - P_r W_{rs}) (\ln P_r - \ln P_s)$$

↓

$$\frac{dH}{dt} = -\frac{1}{2} \sum_r \sum_s (P_r W_{rs} - P_s W_{sr}) (\ln P_r - \ln P_s)$$

In this case we can no longer be sure that the argument of the sum is ≥ 0 . The weaker condition we get

so that $\frac{dH}{dt} \leq 0$ if =

$$\ln P_r - \ln P_s \geq 0 \quad \text{AND} \quad (P_r W_{rs} - P_s W_{sr}) \geq 0$$

↓

$$P_r > P_s$$

↓

$$1 < \frac{P_r}{P_s} > \frac{W_{sr}}{W_{rs}}$$